

Entanglement entropy in quantum spin chains with finite range interaction

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B

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Work in collaboration with Alexander Its and Man Yue Mo

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Consider a one-dimensional quantum spin chain:

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Consider a one-dimensional quantum spin chain:

Look at the ground state $j \quad g \quad ih \quad g \quad j$

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Consider a one-dimensional quantum spin chain:

Look at the ground state $\psi_{g,h}(\sigma_j)$ ($T = 0$: phase transition in the thermodynamic limit.)

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Questions we can ask:

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The problem

Consider a one-dimensional quantum spin chain:

Look at the ground state $\psi_{g,h}(\sigma_j)$ ($T = 0$: phase transition in the thermodynamic limit.)

Questions we can ask:

What is the entropy of the entanglement between A and B as $L \rightarrow \infty$?

What is the correlation between two spins at different sites?

Many others.

If the model is:

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Jin and Korepin (2004):

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Jin and Korepin (2004): Spatial isotropy ($\alpha = 0$), next neighbour interaction and translation invariance (XX model, Toeplitz determinants).

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Keating and Mezzadri (2004): Spatial isotropy ($\beta = 0$), finite range interaction, translation invariance and reflection symmetries (β) averages over the classical compact groups (Toeplitz + Hankel determinants.)

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The problem

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Its, Jin and Korepin (2006): Spatial anisotropy ($\beta \neq 0$), next neighbour interaction and translation invariance. (XY model, block-Toeplitz determinants.)

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Its, Mezzadri and Mo (2007):

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Spin chains and block-Toeplitz determinants

The entropy of the entanglement can be written as

$$S(A) = \lim_{l \rightarrow 0^+} \lim_{L \rightarrow \infty} \frac{1}{L} \log \det D_L(\lambda) \quad (1)$$

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Spin chains and block-Toeplitz determinants

$$g^2(z) = \prod_{j=1}^{2n} z - z_j$$

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$$g^2(z) = \dots^2$$

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Remarks:

The branch cuts of $g(z)$ are the segments

$$i = [2i-1, 2i]; \quad i = 1, \dots, 2n$$

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Remarks:

The branch cuts of $g(z)$ are the segments

$$i = [2i - i; 2i]; \quad i = 1; \dots; 2n$$

$g(z)$ has discontinuities $g_+(z) = g_-(z); \quad z \in i$

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Remarks:

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$$i = [2i-1, 2i]; \quad i = 1, \dots, 2n$$

$g(z)$ has discontinuities $g_+(z) = g_-(z); \quad z \in [2i-1, 2i]$
 $g(z)$ lives on the hyperelliptic curve

$$L : w^2 = \prod_{i=1}^{2n} (z - \lambda_i)$$

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 $g(z)$ lives on the hyperelliptic curve

$$L : w^2 = \prod_{i=1}^{2n} (z - i)$$

The genus of L is $g = 2n - 1$.

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Define $(s) : \mathbb{C}^g \rightarrow \mathbb{C}$ associated to L by

$$(s) := \prod_{n \in \mathbb{Z}^g} e^{i \langle s, n \rangle - \langle n, n \rangle / 2}.$$

is a $g \times g$ symmetric matrix (period matrix) that depends on the Hamiltonian through the branch cuts of L .

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What happens at a phase transition?

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How do we compute such formulae?

The symbol $(z) = \frac{i}{g^{-1}(z)}$ ~~is the Jost admittance $J(z)$ of the Wiener-Hopf~~

$$(z) \neq =$$

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How do we compute such formulae?

The symbol $(z) = \frac{i}{g^{-1}(z)} \frac{g(z)}{i}$ admits the Wiener-Hopf factorization:

$$(z) = U_+(z)U_-(z) = V_-(z)V_+(z);$$

where $U_-(z)$ and $V_-(z)$ are analytic inside/outside the unit circle and

$$U_-(1) = V_-(1) = I;$$

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Theorem (Widom, 1974)

$$\frac{d}{dz} \log D_L(z) = \frac{2}{1-z^2} L + \frac{1}{2} \operatorname{tr}_{S^1} U_+^0$$

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It turns out that

$$\begin{aligned} V(z) &= {}_3U(z)^{-1} {}_3 \\ V_+(z) &= {}_3U_+(z)^{-1} {}_3(2 \quad 1); \quad \in 1 \end{aligned}$$

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Let us define

$$\begin{aligned} S(z) &= U_-(z)Q(z)^{-1}; & |z| < 1; \\ S(z) &= U_+(z)^{-1}Q(z); & |z| > 1; \end{aligned}$$

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In order to compute the entropy of entanglement we need:

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In order to compute the entropy of entanglement we need:

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In order to compute the entropy of entanglement we need:

- 1 to solve the previous RH problem for $S(z)$ in terms of) ~~0564~~

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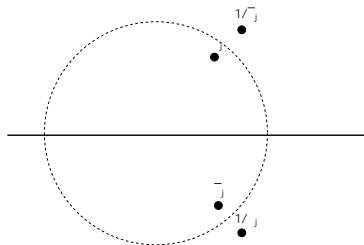
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The critical case (phase transitions)

Pairs of roots of $g(z)$ approach the unit circle.



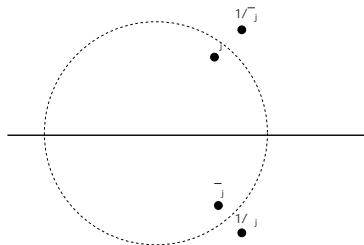
The period matrix in

$$(\mathbb{1}_s) := \prod_{\substack{j \in \mathbb{Z} \\ |j| \leq n}} e^{j \mathbb{1}_n - \mathbb{1}_n - 2i \mathbb{1}_s \mathbb{1}_n};$$

becomes degenerate.

The critical case (phase transitions)

Pairs of roots of $g(z)$ approach the unit circle.



The period matrix \mathcal{X} in

$$(\mathcal{I}^s) := \prod_{n \in 2\mathbb{Z}^g} e^{j \cdot \mathcal{I}^n} \mathcal{I}^n 2i \mathcal{I}^s \mathcal{I}^n ;$$

becomes degenerate.

(\mathcal{I}^s) becomes singular.

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We computed the entropy of the entanglement of the ground state of integrable quantum spin chains with finite range and translation invariant interaction.

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At the core of the computation is the evaluation of block-Toeplitz determinants.

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We computed the entropy of the entanglement of the ground state of integrable quantum spin chains with finite range and translation invariant interaction.

At the core of the computation is the evaluation of block-Toeplitz determinants.

Such determinants are computed by solving a RH problem.

At phase transition we observe logarithmic divergences that generalize previous results.

Summary

We computed the entropy of the entanglement of the ground state of integrable quantum spin chains with finite range and translation invariant interaction.

At the core of the computation is the evaluation of block-Toeplitz determinants.

Such determinants are computed by solving a RH problem.

At phase transition we observe logarithmic divergences that generalize previous results.

AR Its, F Mezzadri and MY Mo. Entanglement entropy in quantum spin chains with finite range interaction.
arXiv: 0708.0161v1.